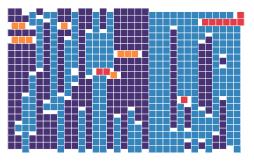
Random high-density packings of  $2 \times 2$  tiles on the square lattice

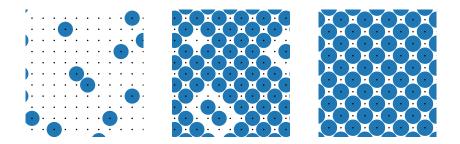
Daniel Hadas, Tel Aviv University Joint work with Ron Peled



HUJI Dynamics Seminar, 16/11/2021

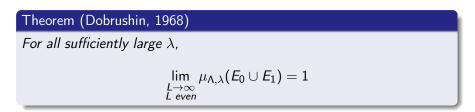
#### The hard-core lattice gas

- Domain:  $\Lambda = (\mathbb{Z}/L\mathbb{Z})^2$  ("a discrete  $L \times L$  torus").
- Configuration: a set  $\sigma \subset \Lambda$  with no two points at distance one.
- Fugacity parameter:  $\lambda > 0$ .
- Probability measure:  $\mu_{\Lambda,\lambda}(\sigma) = \frac{\lambda^{\#\sigma}}{Z_{\Lambda,\lambda}}$  where:
  - $\#\sigma$  is the number of points in  $\sigma$ .
  - $Z_{\Lambda,\lambda}$  is a normalization constant (the partition function).



#### Phase transition in the hard-core model

- Restrict to even *L*. The long-range order is captured by two events:
- $E_0 = \{99\% \text{ of the points in } \sigma \text{ have an even sum of coordinates}\}.$
- $E_1 = \{99\% \text{ of the points in } \sigma \text{ have an odd sum of coordinates}\}.$



- Dobrushin uniqueness condition: model is disordered at small  $\lambda$ .
- Open: Is there a single transition point  $\lambda_c$  from a disordered to an ordered state?

## Continuum hard-core models

- Configuration  $\sigma$  consists of spheres in a domain in  $\mathbb{R}^d$ .
- Sampled with probability proportional to  $\lambda^{\#\sigma}$  (with respect to a suitable Lebesgue measure).
- Major open problems:

Is there an ordered state in dimensions  $d \ge 3$ ? Is the rotational symmetry broken in dimension d = 2? Richthammer (2007): No translational-symmetry breaking in two dimensions.

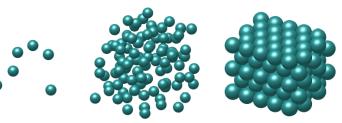


image credit: Ian Jauslin

#### Lattice hard-core models

- Comprehensive study by Mazel–Stuhl–Suhov (2018-19) of hard-core models on Z<sup>2</sup>, triangular and hexagonal lattices with general radius of exclusion.
- Prove long-range order at high fugacities in non-sliding cases.
- Sliding phenomenon:
  - Significant non-uniqueness of maximal density packings due to a sliding degree of freedom.
  - Occurs for a finite number of exclusion radii.
  - Unclear whether these cases still undergo a phase transition.

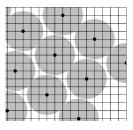
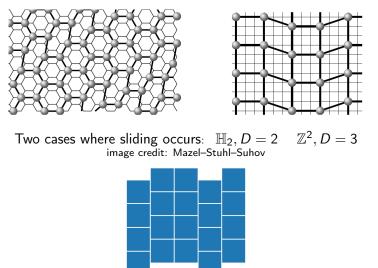


image credit: Izabella Stuhl

#### Lattice hard-core models: Sliding



Sliding in the 2 × 2-hard-square model ( $\mathbb{Z}^2$ , D = 2)

# Hard rod models

- Monomer-Dimer model:
  - Configuration consists of  $2 \times 1$  rods (i.e., a matching).
  - Heilmann–Lieb (1972) famously proved the absence of a phase transition on all graphs.
- Many other models:
  - Onsager (isotropic-nematic transition in liquid crystals, 1949)
  - Heilmann–Lieb (1979) and Jauslin–Lieb (interacting monomer-dimer 2018).
  - Ioffe-Velenik-Zahradník (variable length rods on  $\mathbb{Z}^2$ , 2005)
  - Disertori–Giuliani (long rods on  $\mathbb{Z}^2$ , 2013),
  - Disertori–Giuliani–Jauslin (anisotropic plates in  $\mathbb{R}^3$ , 2020)

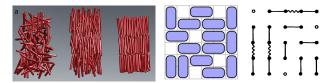
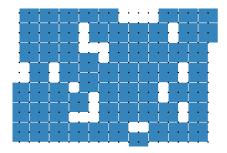


image credits: Zvonimir Dogic (2016), Nicolas Allegra (2015), Heilmann-Lieb (1979)

## The $2 \times 2$ hard squares model

- Domain:  $\Lambda = (\mathbb{Z}/L\mathbb{Z})^2$  ("a discrete  $L \times L$  torus").
- Configuration: a set  $\sigma$  of pairwise disjoint 2  $\times$  2 tiles with centers in A.
- Fugacity parameter:  $\lambda > 0$ .
- Probability of a configuration:  $\mu_{\Lambda,\lambda}(\sigma) = \frac{\lambda^{\#\sigma \frac{L^2}{4}}}{Z_{\Lambda,\lambda}}$  where:
  - $\#\sigma$  is the number of tiles in  $\sigma$ .  $\left(-4\left(\#\sigma - \frac{l^2}{4}\right)$  counts vacant  $1 \times 1$  squares).
  - $Z_{\Lambda,\lambda}$  is a normalization constant (the partition function).



## Main result: Columnar order

- Restrict to even L.
- Each tile has one of four parities: (Even,Even), (Even,Odd), (Odd,Even), (Odd, Odd).
- Let  $E_{\parallel,0}$ , be the "ordering by even columns" event: more than 49% of the tiles have parity (Even, Even), and more than 49% of the tiles have parity (Even, Odd).
- Similarly define  $E_{|,1}, E_{-,0}, E_{-,1}$ .

Theorem (H.–Peled, 2021+)

For all sufficiently large  $\lambda$ ,

$$\lim_{\substack{L\to\infty\\ even}} \mu_{\Lambda,\lambda}(E_{|,0}\cup E_{|,1}\cup E_{-,0}\cup E_{-,1}) = 1$$



an illustration of  $E_{\parallel,0}$ 

#### Theorem (H.–Peled, 2021+)

For sufficiently large  $\lambda$ , the following holds:

The set of doubly-periodic infinite volume Gibbs measures, is a simplex with four vertices (denoted  $\mu_{ver,0}$ ,  $\mu_{ver,1}$ ,  $\mu_{hor,0}$  and  $\mu_{hor,1}$ ).

These four measures are related to each other by translations and rotations.

One of them  $(\mu_{\mathrm{ver},0})$  satisfies the following:

- $\mu_{ver,0}$  is  $(2\mathbb{Z} \times \mathbb{Z})$ -invariant and extremal.
- Solumnar order:  $\mu_{\text{ver},0}(\sigma(0,1)) = \Theta(\lambda^{-1}).$
- Orrelations decay exponentially with distance, for a non-isotropic distance function:

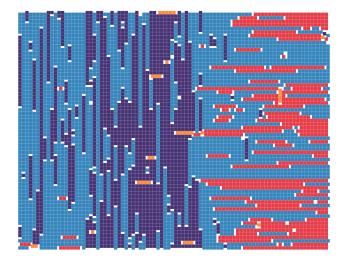
$$d_{ ext{ver}}((x_1,y_1),(x_2,y_2)):=\lambda^{-1/2}|y_2-y_1|+|x_2-x_1|$$



# Proof Ideas

• Note: we only discuss the proof of orientational order.

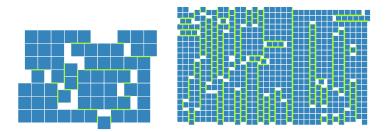
# Interfaces between phases



## Sticks

For a configuration  $\sigma$ , define:

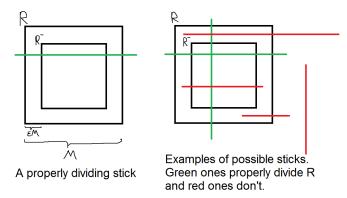
- a stick edge: a segment of length 1, bounding on tiles of different parities.
- a stick: a maximal path of stick edges.



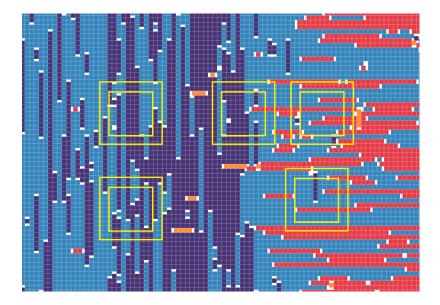
- Sticks cannot intersect. Thus two close long sticks must have same orientation.
- Bound the probability that most sticks are short, by direct calculation.

# Properly divided squares (1/3)

- Let  $\epsilon$  be a small constant.
- Set  $M = M(\lambda)$  (think  $M = \epsilon \lambda^{1/2}$ )
- Let R be a  $M \times M$  square
- define  $R^-$  as  $(1 2\epsilon)M \times (1 2\epsilon)M$  square concentric to R. (assume  $\epsilon M \in \mathbb{Z}$ )
- Say R is properly divided (for  $\sigma$ ) if a stick divides both R and  $R^-$ .

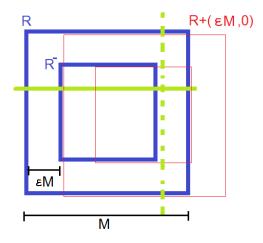


# Properly divided squares (2/3)



# Properly divided squares (3/3)

- If R and  $R + (\epsilon M, 0)$  are properly divided, they are properly divided in same orientation.
- Same for R and  $R + (0, \epsilon M)$ .



#### The main lemma

Let *R* be  $M \times M$  for  $M = \epsilon \lambda^{1/2}$ . Denote by  $E_R$  the event that *R* is **not** properly divided.

#### Lemma

There is  $\epsilon > 0$  such that for all sufficiently large  $\lambda$ ,

 $\mu(E_R) \leq e^{-\epsilon^3 \lambda^{1/2}}$ 

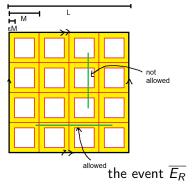
 In fact a multiplicative bound holds. If A is a set of copies of R shifted by vectors in (MZ)<sup>2</sup> then

$$\mu(\bigcap_{R'\in A} E_R) \le e^{-\epsilon^3\lambda^{1/2}|A|}$$

• This allows to prove orientational order with a Peierls argument.

# The disseminated event

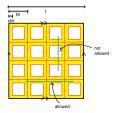
• Define the disseminated version of  $E_R$  to be  $\overline{E_R} := \bigcap_{v \in (M\mathbb{Z})^2/(L\mathbb{Z})^2} E_{R+v}$ .



- The  $2 \times 2$  hard-square model is satisfies reflection positivity.
- Thus the chessboard estimate holds, and implies:

$$\mu(E_R) \leq \left(\mu(\overline{E_R})\right)^{\frac{M^2}{L^2}}$$

#### Bounding the disseminated event



- In  $E_R$  all sticks have length 2*M* at most, except for sticks contained in the yellow regions.
- For simplicity we will discuss bounding the sum over the event  $E_M$  that all sticks have length at most 2M.

$$\mu(E_M) = \frac{\sum_{\sigma \in E_M} \lambda^{\#\sigma - \frac{L^2}{4}}}{Z_{\Lambda,\lambda}} \leq ?$$

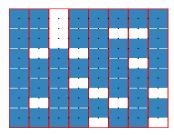
• Bound separately the nominator and denominator.

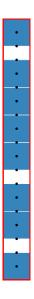
# Lower bound on $Z_{\Lambda,\lambda}$

- Consider a one dimensional system: Λ = {0} × Z/LZ for even L.
- Geometrically, configurations are packings in a 2  $\times$  L rectangle.

• Easy to see that: 
$$Z_{\{0\} imes\{1,2...,L-1\},\lambda} \geq \left(1+\lambda^{-1/2}
ight)^{L/2}$$

• Conclude for the torus:  $Z_{(\mathbb{Z}/L\mathbb{Z})^2,\lambda} \ge (1 + \lambda^{-1/2})^{L^2/4} \approx e^{\frac{1}{4}L^2\lambda^{-1/2}}$ 





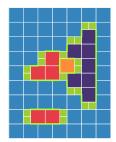
# Bounding $\sum_{\sigma \in E_M} \lambda^{\#\sigma - L^2/4}$

- $\mathcal{H}_M = \{\text{components of vacancies and sticks that may appear in } \sigma \in E_M, \text{ up to translation} \}$
- weight of a component:  $w(c) := \lambda^{-\frac{1}{4}v(c)}$  where v(c) is the number of vacancies in  $c \in \mathcal{H}_M$ .

• Next slide: 
$$\sum_{c\in\mathcal{H}_M}w(c)=C\epsilon\lambda^{-1/2}$$

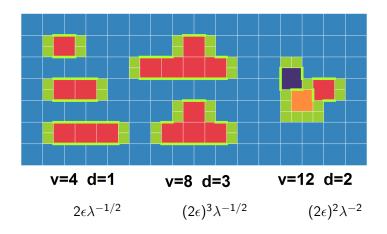
• For each component pick an arbitrary root.

$$\sum_{\sigma \in E_M} \lambda^{\#\sigma - \frac{L^2}{4}} \leq \sum_{\sigma \in E_M} \prod_{v \in \Lambda^2} \begin{cases} w(c) & v \text{ root of } c \\ 1 & o/w \end{cases}$$
$$\leq \left( 1 + \sum_{c \in \mathcal{H}_M} w(c) \right)^{L^2} \leq e^{C\epsilon \lambda^{-1/2} L^2}$$



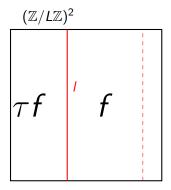
Bounding  $\sum_{c\in\mathcal{H}_M}w(c)$ 

Proof idea: Sum λ<sup>-1/4 v</sup>(2M)<sup>d</sup> = (2ε)<sup>d</sup> · λ<sup>(2d-v)/4</sup> over "components up to the length of sticks" where d counts "degrees of freedom".



# Reflection positivity

- I a vertical line through vertices of  $\boldsymbol{\Lambda}$
- I and its opposite divide  $\Lambda$  to two rectangles  $R_0, R_1$ .
- $\tau$  is the reflection through  $\mathit{I}$
- $\tau$  exchanges  $R_0$  with  $R_1$ .
- let f be  $R_0$ -local function.
- Conditioned on the restriction to *I* and its opposite,  $\mu(f) = \mu(\tau f)$  and  $\mu(f \cdot \tau f) = \mu(f) \cdot \mu(\tau f)$  thus, Reflection positivity:  $\mu(f \cdot \tau f) \ge 0$



### Reflection positivity

- for  $R_0$ -local f, g, define  $\langle f, g \rangle = \mu(f \cdot \tau g)$
- Reflection positivity:  $<\cdot,\cdot>$  is a non-negative bilinear form.
- Thus the Cauchy-Schwarz inequality holds:  $< f, g > \le \sqrt{< f, f > < g, g >}$
- Example:  $\mu(f) \leq \sqrt{\mu(f \cdot \tau_{s_1} f)} \leq \sqrt[4]{\mu(f \cdot \tau_{s_1} f \cdot \tau_{s_2} f \cdot \tau_{s_1} \tau_{s_2} f)}$

	S <sub>1</sub>	
<b>S</b> <sub>2</sub>	f	

#### The chessboard estimate

- Let R be a rectangle.
- assume 2Width(R), 2Height(R) divide L.
- Let T = T<sup>R</sup><sub>Λ</sub> be the isometries generated by reflections in the sides of R.
- For each  $\tau \in T$ , let  $f_{\tau}$  be *R*-local. (Then  $\tau f_{\tau}$  is  $\tau R$ -local)
- Define a norm:

$$\mathfrak{z}^{R}_{\Lambda}(f) \coloneqq \left[ \mu \left( \prod_{\tau \in T} \tau f \right) \right]^{1/\#T}$$

Then

$$\mu\left(\prod_{\tau\in T}\tau f_{\tau}\right)\leq\prod_{\tau\in T}\mathfrak{z}_{\Lambda}^{R}(f_{\tau})$$

f	ţ	f	4
f	f	f	f
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- A similar result is expected for  $k \times k$  tiles, however, a different proof is needed since reflection positivity does not apply.
- What happens for 2 × 2 × 2 cubes? We conjecture the existance of exactly 12 phases (of columnar order) at high fugacity.
- What happens for 1 × k rods? At intermediate fugacity, a nematic phase was proved using cluster expansions (Disertori and Giuliani, 2013). What happens at high fugacity?

you for listening!